

From forward modeling of acoustic wave equation towards to reverse time migration (rtm)

D.A.Gutiérrez

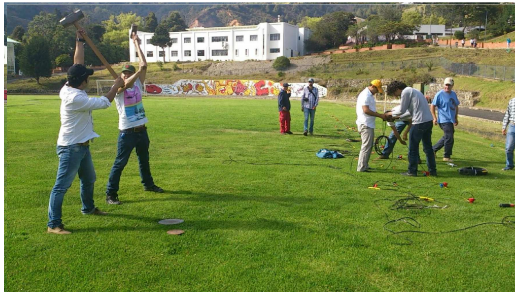
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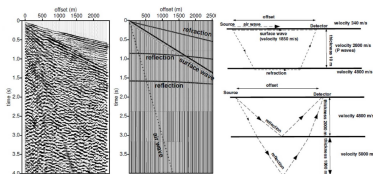
Adquisition of seismic data



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Overview to seismic processing (Schuster, 2010 (book))

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5. Repeat steps 3-4 for all midpoint gathers to give the seismic section

- ▶ Migration (Schuster, 2010)
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- ▶ RTM (Beyreht et al., 2002)

$$P(x, z, t) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \hat{P}(x, z, \omega) e^{-i\omega t} d\omega$$

Richardson et al. (1980) (Schuster et al., 1977)

Wang and Claessens (2000)



- ▶ Migration (Schuster, 2010)
- ▶ Inversion (Schuster, 2010)
- ▶ RTM (Baysal *et. al.*, 1983)

$$[P(x, z, T + \Delta t) - P(x, z, T - \Delta t)]/2\Delta t = \dot{P}(x, z, T)$$

Including reflectivity model (Lisowski *et. al.*, 1976)

Using Casprini (Casprini, 1982)

- ▶ Migration (Schuster, 2010)
Is defined as the process which takes the seismic section $d(x,z,t)$ and moves the reflection events back to their origin at the interfaces

$$\mathbf{d} = \mathbf{L}(\mathbf{m})$$

- ▶ Inversion (Schuster, 2010)
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Exploding reflector model (Loewentgal *et. al.*, 1976)
Imaging Condition (Clearbout, 1985)

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Full wave inversion can be described as an iterative sequence of migrations, where the data residuals updated and migrated at each iteration to give the new model update, see

$$\mathbf{m}^{(k+1)} = \mathbf{m}^{(k)} + [\mathbf{L}_{(k)}^T \mathbf{L}_{(k)}]^{-1} \mathbf{L}_{(k)}^T \mathbf{d}^{(k)}$$

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An approximation to a stacked section can be obtained in a single experiment by replacing the subsurface with a medium containing half the actual velocities in the earth, and by initiating explosive sources at time zero on all the reflecting boundaries **Imaging Condition** (Clearbout, 1985)

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$$\begin{aligned} \mathbf{P}(\vec{x}, t) &= \mathbf{P}_s(\vec{x}, t) \star \mathbf{P}_r(\vec{x}, t) \\ \mathbf{R}(\vec{x}) &= \mathbf{P}(\vec{x}, t = 0) \sim \mathbf{m} \end{aligned}$$

Acoustic limit of the elastic wave propagation

The consequences. . .

The consequences of this acoustic approximation include the restrictions to isotropic source radiation patterns and absence of the Rayleigh waves and P-to-S conversions. The acoustic approximation is nevertheless, justifiable when data analysis is restricted to the first-arriving P waves and when the seismic sources radiate little S wave energy (e.g. explosions) (Fichtner, 2011)^a

^aPage 14, Andreas Fichtner, "Full Seismic Waveform Modelling and Inversion", Advances in Geophysical and Environmental Mechanics and Mathematics, Springer(2011)

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$$\begin{aligned}\rho(\vec{x})\ddot{\vec{u}}(\vec{x}, t) - \nabla \cdot [\mathbb{C}(\vec{x}) : \nabla \vec{u}(\vec{x}, t)] &= \vec{f}(\vec{x}, t) \\ \sigma(\vec{x}, t) &= \mathbb{C}(\vec{x}) : \nabla \vec{u}(\vec{x}, t)\end{aligned}$$

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- ▶ Acoustic Limit

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- ▶ Acoustic Limit

$$\begin{aligned}\rho(\vec{x})\partial_t \vec{u}(\vec{x}, t) &= -\nabla p(\vec{x}, t) \\ \frac{1}{\rho(\vec{x})c(\vec{x})^2} \partial_t p(\vec{x}, t) &= -\nabla \cdot \vec{u}(\vec{x}, t)\end{aligned}$$

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Diffusion Equations 2D ($\rho(\vec{x}) = 1, c(\vec{x}) = 1$)

$$\partial_t p(\vec{x}, t) + \partial_x u(\vec{x}, t) + \partial_z v(\vec{x}, t) = 0$$

$$\partial_t u(\vec{x}, t) + \partial_x p(\vec{x}, t) = 0$$

$$\partial_t v(\vec{x}, t) + \partial_z p(\vec{x}, t) = 0$$

- ▶ Pseudospectral Method (PS) via Finite Differences (FD) (t, \vec{k})
PS-FD System

$$\bar{\mathbf{P}}(t + \Delta t) - \bar{\mathbf{P}}(t) = -\Delta t(\mathcal{F}^{-1}[ik_x \mathcal{F}(u)] + \mathcal{F}^{-1}[ik_z \mathcal{F}(v)])$$

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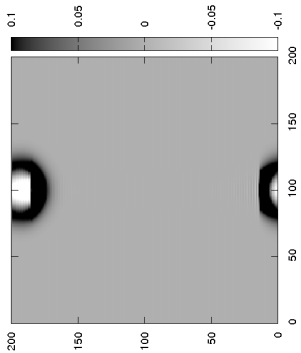
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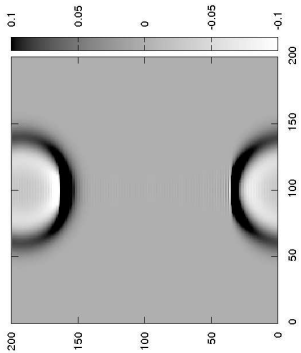
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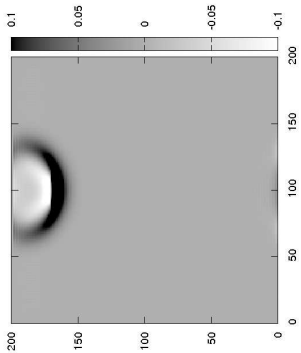
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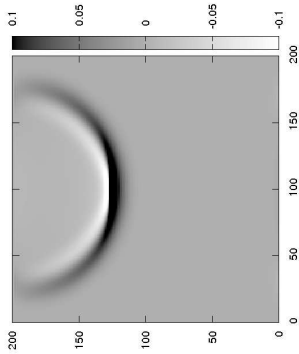
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- ▶ Results. . .









Another works

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1. Modeling in acoustic variable-density media by Fourier finite differences (Xiaolei Son, 2012)

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8. There is more ...

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2. From the elastic wave equation we derive stress-velocity diffusion system equations on which we take the acoustic limit restricting our data analysis
3. The Pseudospectral methods via finite differences are implemented in acoustic media as a trial platform and it offers interesting possibilities in inhomogeneous media